

# Sampling from spherical spin glasses in total variation via algorithmic stochastic localization

Brice Huang (MIT)

JSM: Advances in the theory of modern sampling algorithms  
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Joint work with Andrea Montanari and Huy Tuan Pham (Stanford)

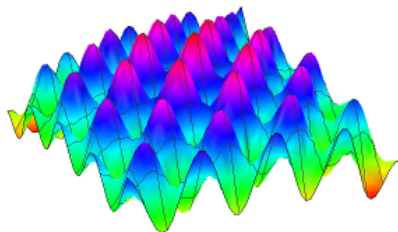
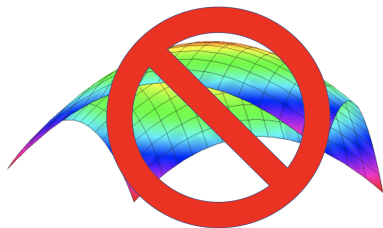


# Goal of this talk

Sample from a **high-dimensional** measure

$$\sigma \sim \mu(d\sigma) \equiv \frac{1}{Z} e^{\beta H_N(\sigma)} d\sigma, \quad \mu \in \mathcal{P}(\mathbb{R}^N)$$

where  $H_N$  **non-concave** and **highly multimodal**



# Mean-field spin glasses

- Random polynomial over  $\{\pm 1\}^N$  or  $S_N = \{\boldsymbol{\sigma} \in \mathbb{R}^N : \|\boldsymbol{\sigma}\| = \sqrt{N}\}$
- Sherrington-Kirkpatrick model:

$$H_N(\boldsymbol{\sigma}) = \frac{1}{\sqrt{N}} \sum_{i,j=1}^N g_{i,j} \sigma_i \sigma_j \quad g_{i,j} \stackrel{i.i.d.}{\sim} \mathcal{N}(0, 1)$$

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- Mixed  $p$ -spin model:

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- $\mu_\beta(d\sigma) \equiv \frac{1}{Z} e^{\beta H_N(\sigma)} d\sigma$  Gibbs measure at inverse temperature  $\beta$

# Connections

Spin glasses are prototypes for disordered, random probability measures:

- Posteriors in high-dimensional Bayesian inference
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Estimate  $\mathbf{x}_0 \sim \text{unif}(\{\pm 1\}^N)$  from noisy observation  $\mathbf{A} = \lambda \mathbf{x}_0^{\otimes 2} + \mathbf{W}$

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Posterior is precisely

$$p(\mathbf{x}|\mathbf{A}) = \frac{1}{Z} \exp(\lambda(\mathbf{A}\mathbf{x}, \mathbf{x}))$$

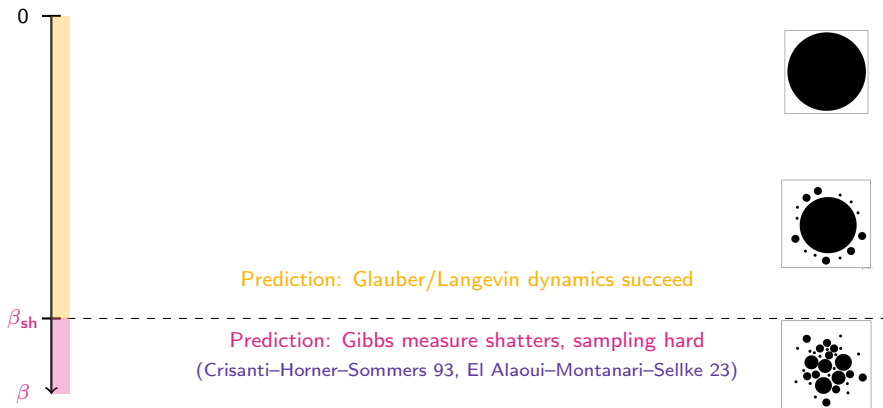
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For sampling from Gibbs measure  $\mu_\beta(d\sigma) \equiv \frac{1}{Z} e^{\beta H_N(\sigma)} d\sigma$ :



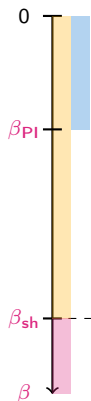
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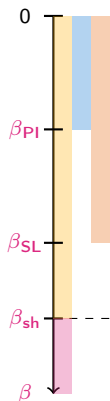


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Algorithmic stochastic localization samples  
with small **Wasserstein** error  
(El Alaoui–Montanari–Sellke 22, Celentano 24)



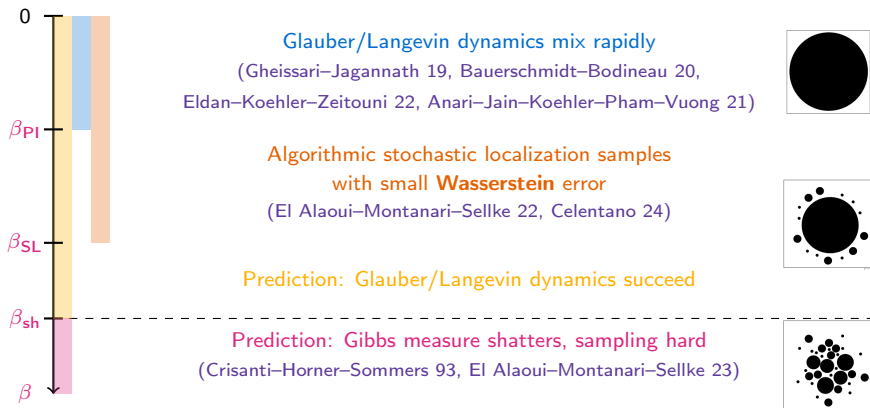
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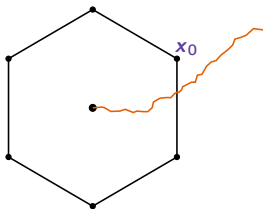
This work: in **spherical models**,  $\beta < \beta_{SL}$ , alg-SL samples in **total variation**

# Stochastic localization

Given  $\mu \in \mathcal{P}(\mathbb{R}^N)$ , consider Brownian motion with **unknown drift**

$$y_t = tx_0 + B_t \quad \sim \mathcal{N}(tx_0, tI_N)$$

Here  $x_0 \sim \mu$  independent of  $B_t$ , and **only  $y_t$  observed**

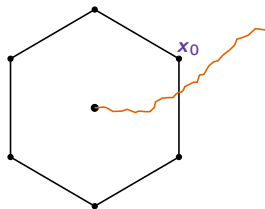


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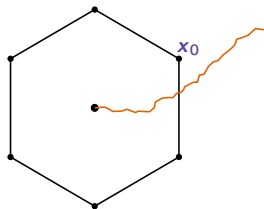


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Reparametrization of “backward process” in denoising diffusions

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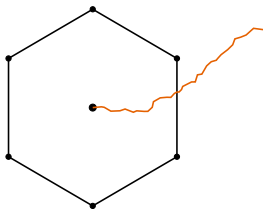
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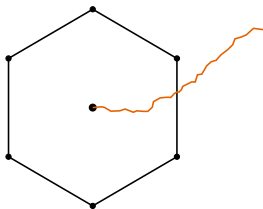
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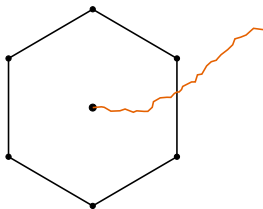
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AMS22: estimate  $\mathbf{m}_t$  by approximate message passing (AMP) iteration

# Results

Recall  $H_N(\boldsymbol{\sigma}) = \sum_p \frac{\gamma_p}{N^{(p-1)/2}} \langle \mathbf{G}^{(p)}, \boldsymbol{\sigma}^{\otimes p} \rangle$ . **Mixture function:**  $\xi(q) = \sum_{p \geq 2} \gamma_p^2 q^p$

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For pure  $p$ -spin model  $\xi(q) = \beta^2 q^p$ ,

$$\frac{\beta_{\text{SL}}}{\beta_{\text{sh}}} \in \left[ \frac{\sqrt{e}}{2} \approx 0.824, 1 \right]$$

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- 2 Sample from  $\mu_T = \mathcal{L}(\mathbf{x}_0 | \mathbf{y}_T)$ , which is **log-concave** for large  $T = O(1)$

## Improved mean estimation

$\mu_t = \mathcal{L}(\mathbf{x}_0 | \mathbf{y}_t)$  is Gibbs measure of **tilted** model:  $\mu_t(d\sigma) \equiv \frac{1}{Z} e^{H_{N,t}(\sigma)} d\sigma$  for

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Estimator of AMS22: solution  $\tilde{\mathbf{m}}_t = \mathbf{m}^{\text{TAP}}$  to **TAP equation**

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Our analysis:

- Actually,  $\mathbb{E} \|\mathbf{m}^{\text{TAP}} - \text{mean}(\mu_t)\|^2 = O(1)$
- Compute correction  $\Delta$ , so  $\hat{\mathbf{m}}_t = \mathbf{m}^{\text{TAP}} + \Delta$  has error  $o(1)$

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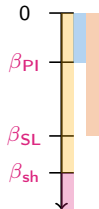
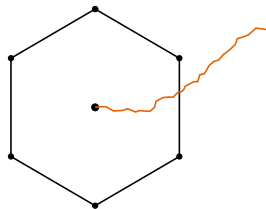
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