

Sampling from spherical spin glasses: diffusions and simulated annealing

Brice Huang (Stanford → Yale)

Workshop on High-Dimensional Learning Dynamics

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Stanford | Science Fellows



Thanks to wonderful collaborators



Andrea
Montanari



Huy Tuan
Pham



Sidhanth
Mohanty



Amit
Rajaraman



David X. Wu

Outline of talk:

Introduction and motivating problems

Results: sampling from spin glasses

Algorithmic stochastic localization

Sampling guarantee for simulated annealing

Conclusion

High-dimensional sampling

Motivating problem

Sample from a high-dimensional probability measure

$$\mu(d\mathbf{x}) \propto e^{V(\mathbf{x})} d\mathbf{x}$$

with a polynomial-time algorithm.

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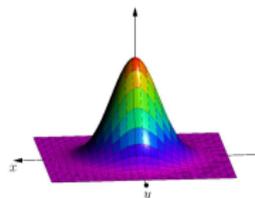
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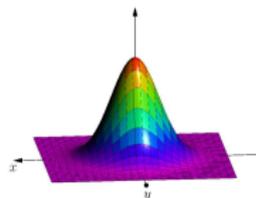
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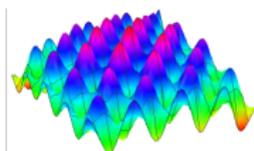
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μ non-log-concave, highly multimodal

Setting of this talk

Model: mean field spin glass

E.g. random cubic polynomial on spherical domain $S_N = \sqrt{N} \cdot \mathbb{S}^{N-1}$:

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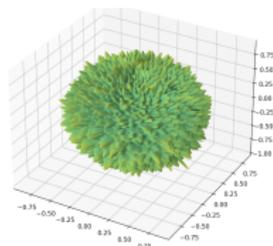
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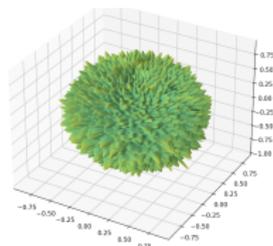
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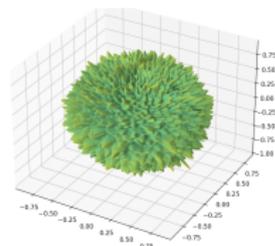
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- linear combination of degrees (**mixed p -spin model**)

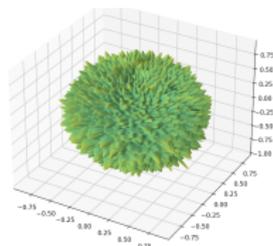
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- linear combination of degrees (**mixed p -spin model**)
- models on cube $\Sigma_N = \{\pm 1\}^N$, e.g. **Sherrington–Kirkpatrick model (1975)** is degree 2

Applications of spin glasses

Prototype of disordered probability measures / objective functions:

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- Random **constraint satisfaction** (MaxCut, MaxSAT, q -coloring):

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C_* defined by a spin glass. (Dembo Montanari Sen 17, Panchenko 18)

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- **Neural network** loss landscapes (Choromanska Henaff Mathieu Ben Arous LeCun 15)
- Posteriors in **high-dimensional Bayesian inference**, e.g. tensor PCA, generalized linear models (Ben Arous Mei Montanari Nica 17, Barbier Krzakala Macris Miolane Zdeborová 18, Fan Mei Montanari 21)

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Problem restatement

We are given a (random instance of) spin glass Hamiltonian, e.g.

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Goal: devise randomized poly-time algorithm outputting \mathbf{x}_{alg} such that

$$\text{Law}(\mathbf{x}_{\text{alg}}) \approx \mu_\beta \quad \text{whp over } H$$

For which β is this possible?

Langevin dynamics

Langevin dynamics

Canonical Markov chain for continuous distributions:

$$d\mathbf{X}_t = \beta \nabla H(\mathbf{X}_t) dt + \sqrt{2} d\mathbf{B}_t \quad \mathbf{B}_t \text{ std } \mathbb{R}^N\text{-Brownian motion}$$

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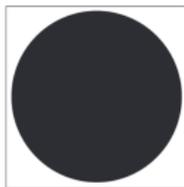
Implement by discretizing time (Chewi Erdođlu Li Shen Zhang 24, ...) but we won't worry about this.

Worst-case vs random initialization

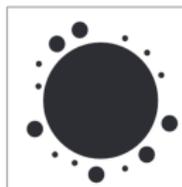
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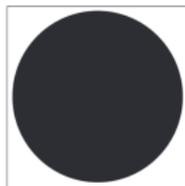
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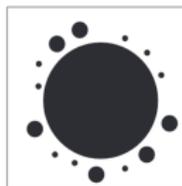
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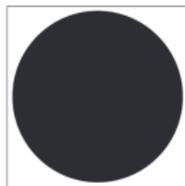


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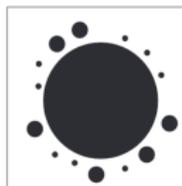
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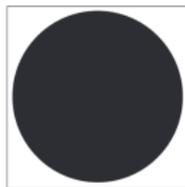
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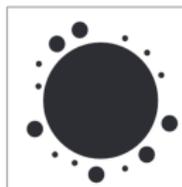
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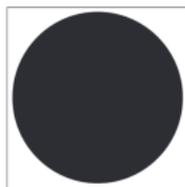
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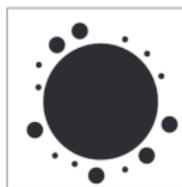
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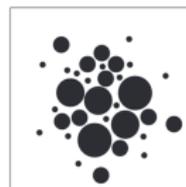
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μ well-connected except $o(1)$
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μ **shattered**

Fails to mix from any
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Results (informal)

Picture conjectured by physics:



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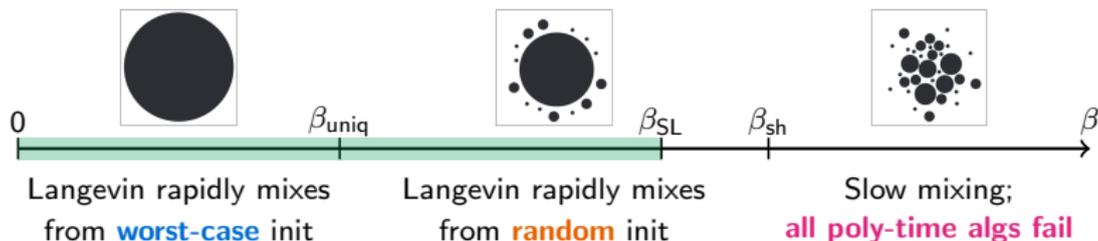


Key physics conjecture

Poly-time algorithms can sample for $\beta < \beta_{\text{sh}}$, and no further. Both sides are open.

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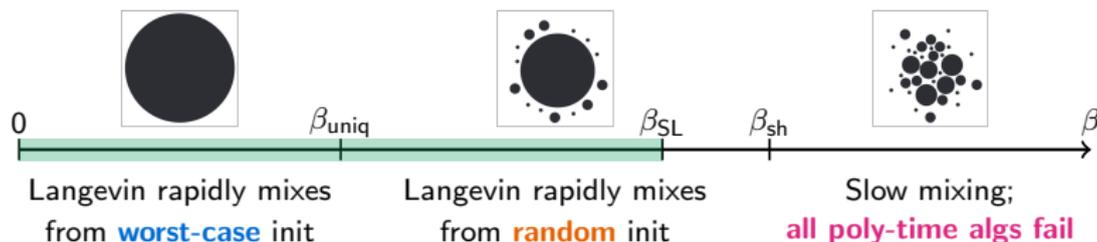
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Main results

Two algorithms succeed up to intermediate β_{SL} : **simulated annealing** and **algorithmic stochastic localization**.

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Main results

Two algorithms succeed up to intermediate β_{SL} :
simulated annealing and **algorithmic stochastic localization**.

For **pure** p -spin model, $\frac{\beta_{\text{SL}}}{\beta_{\text{sh}}} = \frac{\sqrt{e}}{2} + O\left(\frac{1}{p}\right) \approx 0.824$.

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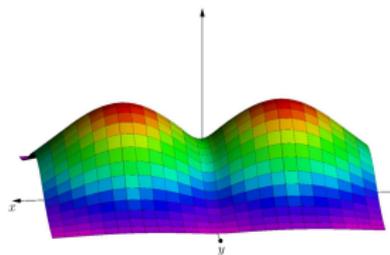


Figure: $\ell/L = 0.1$

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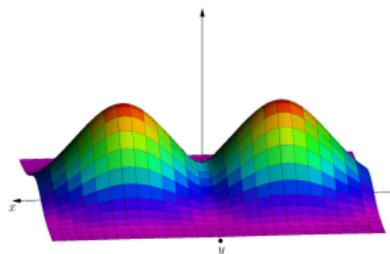


Figure: $\ell/L = 0.2$

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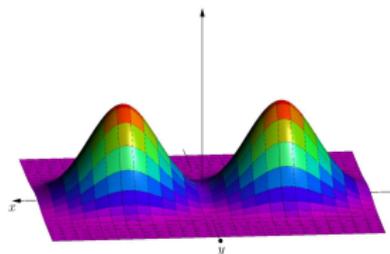


Figure: $\ell/L = 0.4$

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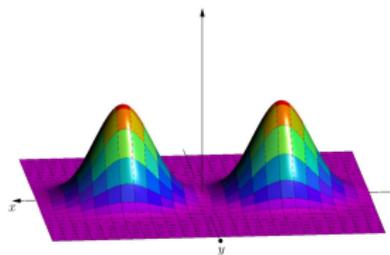


Figure: $\ell/L = 0.6$

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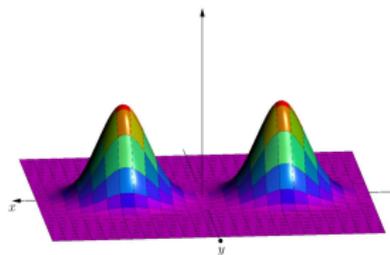


Figure: $\ell/L = 0.8$

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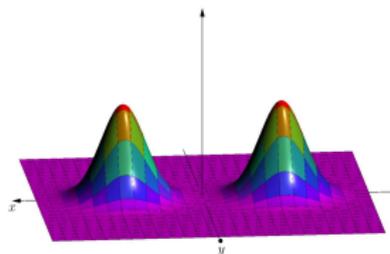


Figure: $\ell/L = 1.0$

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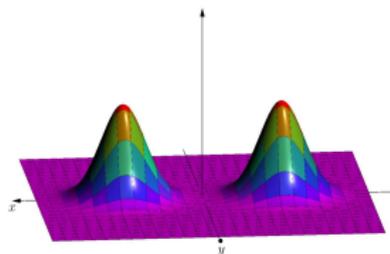


Figure: $\ell/L = 1.0$

Natural and long-studied algorithm for sampling / optimization
(Kirkpatrick Gelatt Vecchi 83, ...) but rigorous guarantees scarce.

Stochastic localization

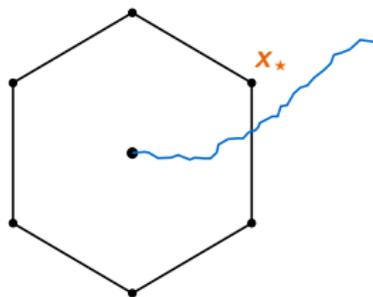
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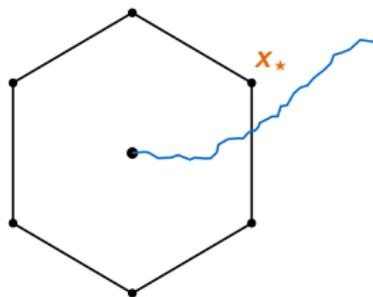
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$$\mu^t = \text{Law}(\mathbf{x}_\star | (\mathbf{y}_s)_{s \leq t})$$

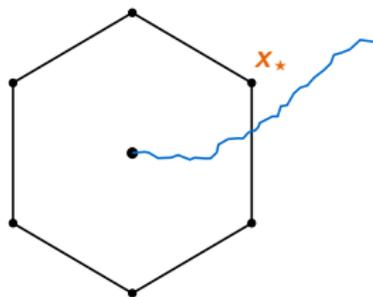
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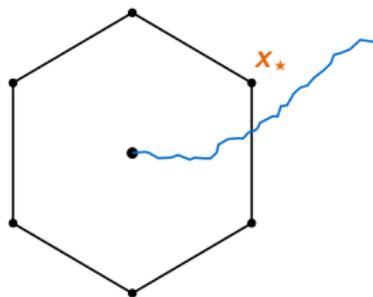
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SL process is the **path of measures** $(\mu^t)_{t \geq 0}$. Localizes to \mathbf{x}_\star as $t \rightarrow \infty$.

Two usages of stochastic localization

SL as a sampling algorithm

Simulate SL process $(\mu^t)_{t \geq 0}$ in distribution (without knowing \mathbf{x}_* !) until $\mu^\infty = \delta_{\mathbf{x}}$ to obtain sample $\mathbf{x} \sim \mu$.

Introduced in [El Alaoui Montanari Sellke 22+23](#) to sample from Sherrington–Kirkpatrick model. Equivalent to denoising diffusions.

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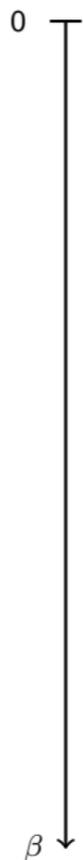
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SL as a proof technique

Control of SL process $(\mu^t)_{t \geq 0} \Rightarrow$ functional ineqs for $\mu \Rightarrow$ rapid mixing of Langevin dynamics on μ .

Originated in convex geometry, breakthrough applications to KLS conjecture ([Eldan 13](#), [Lee Vempala 17](#), [Chen 21](#), [Klartag Lehec 22+25](#))

Results: sampling from spin glass $\mu_\beta(\boldsymbol{\sigma}) \propto e^{\beta H(\boldsymbol{\sigma})}$



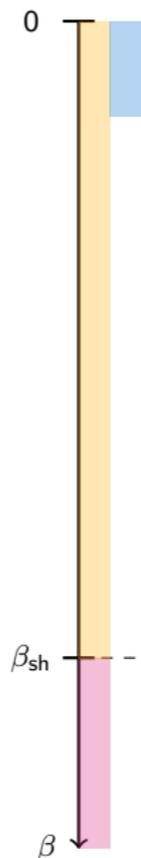
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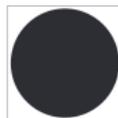


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Langevin mixes rapidly from worst-case init

(Bauerschmidt Bodineau 17, Gheissari Jagannath 19, Eldan Koehler Zeitouni 22, Anari Jain Pham Koehler Vuong 24, AKV 24)



Prediction: Langevin succeeds from random init

Gibbs measure shatters; sampling predicted hard.

Rigorous hardness for **stable** samplers.

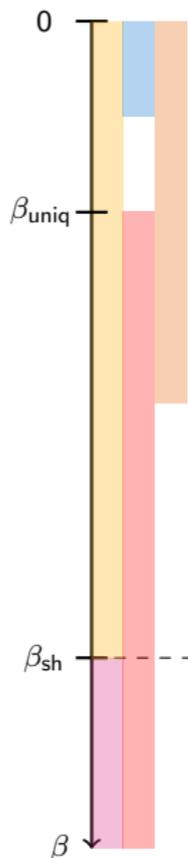
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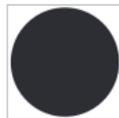


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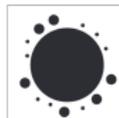
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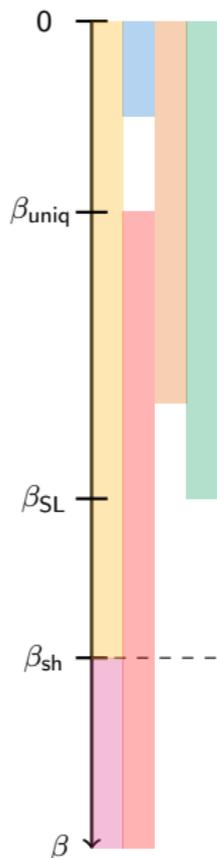
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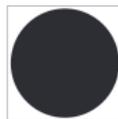
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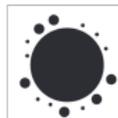


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Result 1: Algorithmic SL succeeds in **total variation** error
+ matching hardness result (H Montanari Pham 24)



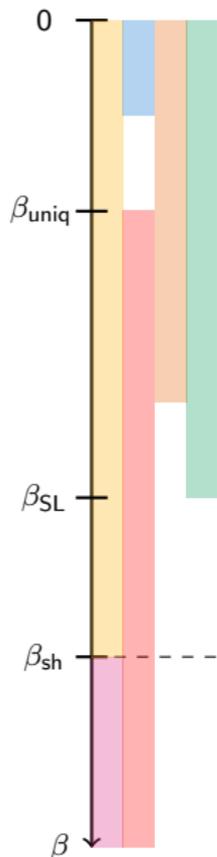
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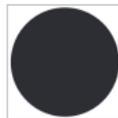


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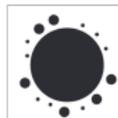


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Result 2: Simulated annealing succeeds (in TV error) (H Mohanty Rajaraman Wu 24)



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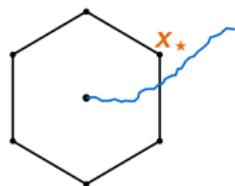
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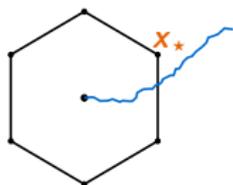
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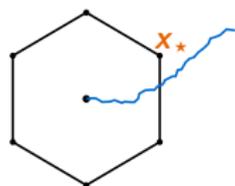
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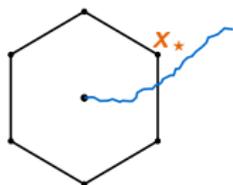
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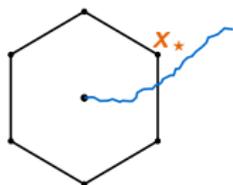
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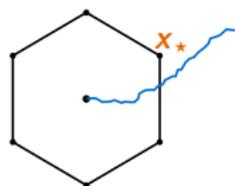
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Upshot: reduces sampling to posterior mean estimation!

Estimate $\mathbf{m}_t \Rightarrow$ simulate $(\mathbf{y}_t)_{t \geq 0} \Rightarrow$ sample from μ

Form of the posterior μ^t

Recall $\mathbf{y}_t = t\mathbf{x}_* + \mathbf{B}_t$. If we observe $\mathbf{y}_t = \mathbf{y}$:

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That is, μ^t is μ with **randomly evolving exponential tilt** (\leftrightarrow external field)

$$\mu^t(d\mathbf{x}) \propto \exp\left(\beta H(\mathbf{x}) + \langle \mathbf{y}_t, \mathbf{x} \rangle\right)$$

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El Alaoui Montanari Sellke 22: estimator \tilde{m}_t for $m_t = \text{mean}(\mu^t)$ by **approximate message passing** (Bolthausen 14, Bayati Montanari 11)

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Main contribution of (H Montanari Pham 24)

For $\beta < \beta_{\text{SL}}$, we develop improved estimator \hat{m}_t with accuracy

$$\mathbb{E}[\|\hat{m}_t - m_t\|^2] = o(1)$$

This implies $d_{\text{TV}}(\mu_{\text{alg}}, \mu_{\beta}) = o(1)$.

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Poincaré inequalities and rapid mixing

For $\mu \in \mathcal{P}(S_N)$, test function $f : S_N \rightarrow \mathbb{R}$, define **Dirichlet form**

$$\mathcal{E}_\mu(f, f) = \mathbb{E}_\mu \left[\|\nabla_{\text{sp}} f\|^2 \right]$$

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μ satisfies a Poincaré inequality with constant $C > 0$ if for all f ,

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❗ Great if this works, but too much to hope for if metastable states ...

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Implies “almost” exponential contraction of χ^2 to μ :

$$\chi^2(P_t \nu \parallel \mu) \leq e^{-Ct} \chi^2(\nu \parallel \mu) + \delta \left\| \frac{d\nu}{d\mu} \right\|_\infty^2$$

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($\delta \|f\|_\infty^2$ accounts for metastable states. Think $\delta = e^{-N^c}$)

Implies “almost” exponential contraction of χ^2 to μ :

$$\chi^2(P_t \nu \parallel \mu) \leq e^{-Ct} \chi^2(\nu \parallel \mu) + \delta \left\| \frac{d\nu}{d\mu} \right\|_\infty^2$$

Consequence: rapid mixing from **warm starts** ν with $\|d\nu/d\mu\|_\infty \ll 1/\sqrt{\delta}$

WPI \Rightarrow simulated annealing succeeds

Consequence of (C, δ) -weak Poincaré inequality

Langevin dynamics rapidly mix to μ from **warm starts** $\left\| \frac{d\nu}{d\mu} \right\|_{\infty} \ll \frac{1}{\sqrt{\delta}}$

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Simulated annealing

Input: β, H . Goal: sample from $\mu_{\beta}(\mathbf{x}) \propto e^{\beta H(\mathbf{x})}$

In stage $\ell = 0, \dots, L$:

- Set $\beta_{\ell} = \beta \cdot \ell/L$. Interpolates $\beta_0 = 0$ to $\beta_L = \beta$
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Related “staging warm starts” idea in [El Alaoui Eldan Gheissari Piana 23](#)

Proving a weak Poincaré inequality

For $\mu \in \mathcal{P}(\mathbb{R}^N)$, $\mathbf{h} \in \mathbb{R}^N$, define **exponential tilt** $\mu^{\mathbf{h}}(\mathrm{d}\mathbf{x}) \propto e^{\langle \mathbf{h}, \mathbf{x} \rangle} \mu(\mathrm{d}\mathbf{x})$

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Message of localization schemes framework (Chen Eldan 22, ...)

μ satisfies a $O(1)$ -Poincaré inequality if

$$\|\mathrm{cov}(\mu^{\mathbf{h}})\|_{\mathrm{op}} = O(1) \quad \text{for all } \mathbf{h} \in \mathbb{R}^N$$

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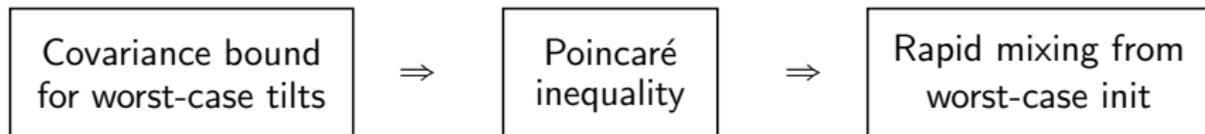
More precisely: \mathbf{h} is the SL tilt \mathbf{y}_t . That is, if

$$\mathbb{P}\left(\sup_{t \geq 0} \|\mathrm{cov}(\mu^{\mathbf{y}_t})\|_{\mathrm{op}} = O(1)\right) \geq 1 - \delta$$

then μ satisfies a $(O(1), \delta)$ -weak Poincaré inequality

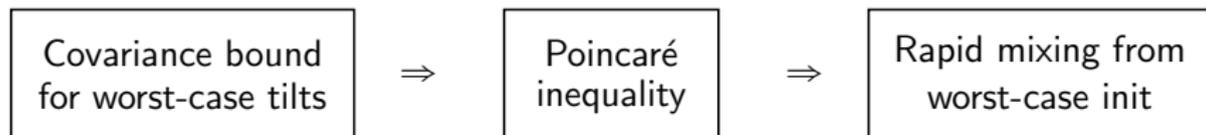
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Approach of [Chen Eldan 22](#):

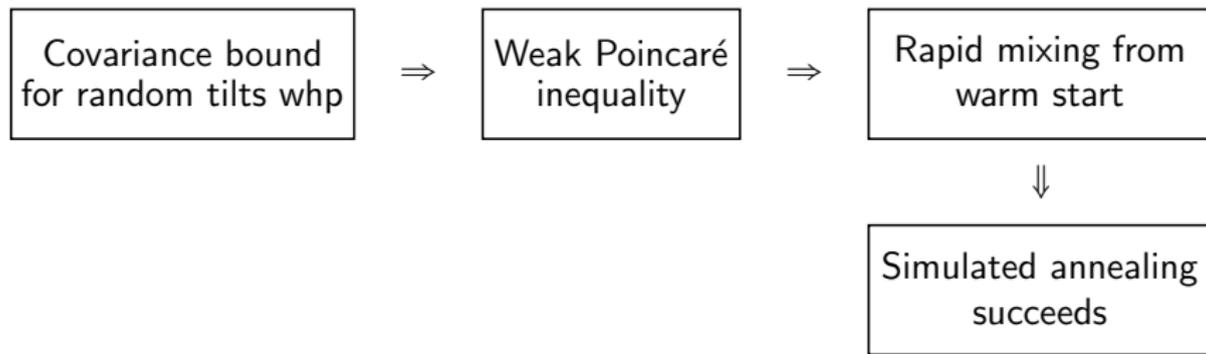


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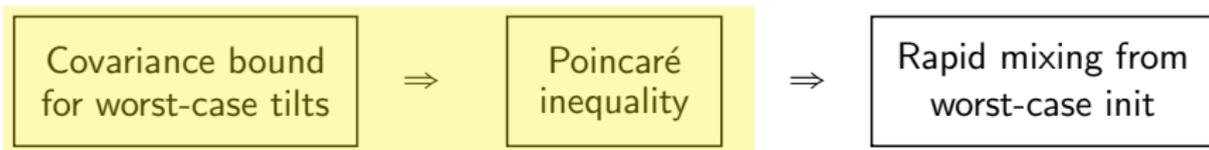


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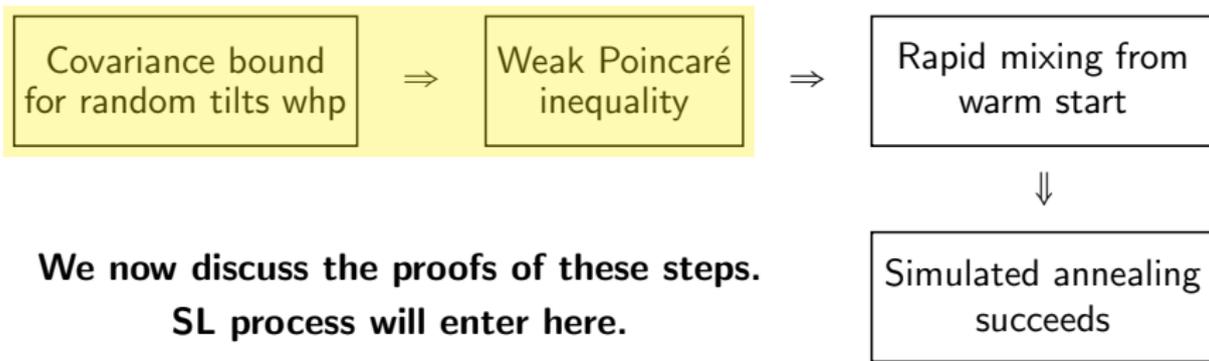


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Approach of [H Mohanty Rajaraman Wu 24](#):



**We now discuss the proofs of these steps.
SL process will enter here.**

Worst-case tilts covariance bound \Rightarrow Poincaré inequality

Let $\mu \in \mathcal{P}(\mathbb{R}^N)$ be target measure,

want to show (for all f)

$$O(1) \cdot \text{Var}_{\mu}(f) \leq \mathcal{E}_{\mu}(f, f)$$

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\forall Supermartingality of \mathcal{E}_{μ^t}

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localized measures have PI

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Main idea of (Chen Eldan 22)

Stochastically differentiating $\text{Var}_{\mu^t}(f)$ yields

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Same strategy, but with a stopping time. Let $T, K = O(1)$ large, and

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Retracing the same steps then proves the WPI.

Outline of talk:

Introduction and motivating problems

Results: sampling from spin glasses

Algorithmic stochastic localization

Sampling guarantee for simulated annealing

Conclusion

Recap: SL and reduced objectives

Goal: sample from Gibbs measure $\mu_\beta(\mathbf{d}\mathbf{x}) \propto \exp(\beta H(\mathbf{x}))$

SL: nature generates secret $\mathbf{x}_\star \sim \mu$, we observe $\mathbf{y}_t = t\mathbf{x}_\star + \mathbf{B}_t$. Posterior:

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Most of the technical content of our papers is to establish these inputs.

One proof idea: re-centering measure with external field

Time t measure μ_{β}^t has external field:

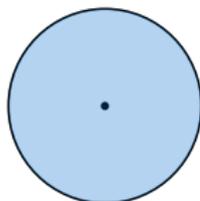
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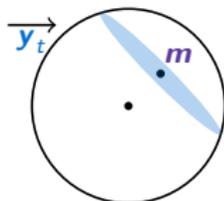
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Gibbs measure concentrates near random codimension-1 band



no field: μ_{β}^t has mean $\mathbf{0}$ and spread “in all directions”



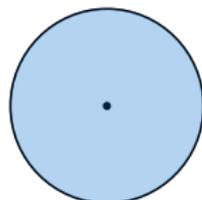
with field: μ_{β}^t concentrates on band with **random** center

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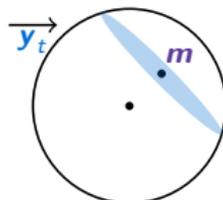
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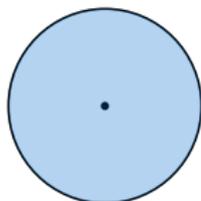
Reduction: restriction to band “looks like” no-field model

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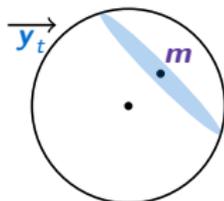
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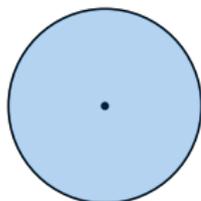
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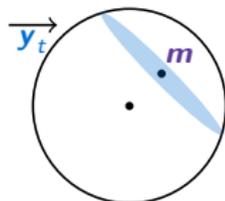
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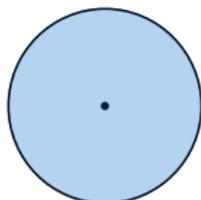
$$\nabla H(m) + \mathbf{y}_t = \left(\frac{1}{1 - \|m\|^2/N} + \left(1 - \frac{\|m\|^2}{N}\right) \xi''\left(\frac{\|m\|^2}{N}\right) \right) m$$

One proof idea: re-centering measure with external field

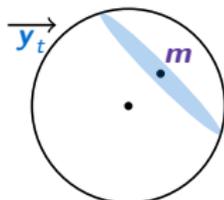
Time t measure μ_β^t has external field:

$$\mu_\beta^t(d\mathbf{x}) \propto \exp\left(\beta H(\mathbf{x}) + \langle \mathbf{y}_t, \mathbf{x} \rangle\right)$$

Gibbs measure concentrates near random codimension-1 band



no field: μ_β^t has mean $\mathbf{0}$ and spread “in all directions”



with field: μ_β^t concentrates on band with **random** center

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Mean estimator: find \mathbf{m} with AMP + GD, plus nontrivial correction

Hardness results

Consider generalization of SL: observe $(\mathbf{x}_*, \mathbf{x}_*^{\otimes 2}, \mathbf{x}_*^{\otimes 3}, \dots)$ through independent gaussian channels, with varying noise levels.

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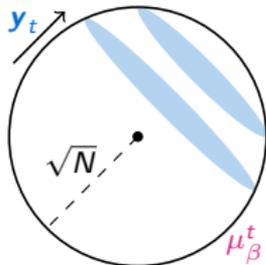
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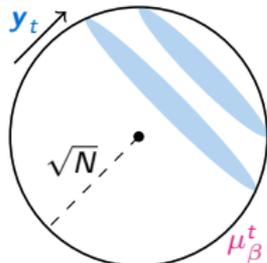
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This also breaks WPI proof: $\|\text{cov}(\mu_\beta^t)\|_{\text{op}} \gg 1$ when μ_β^t straddles two bands



Conclusion

We show two algorithms sample from a spin glass Gibbs measure μ_β in total variation, for all $\beta < \beta_{\text{SL}}$

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Thank you!